Caspian Journal of Applied Mathematics, Ecology and Economics V. 6, No 2, 2018, December ISSN 1560-4055

Weak Solvability of the First Boundary Value Problem for a Class of Parabolic Equations with Discontinuous Coefficients in Paraboloid Type Domains

N.J. Jafarov

Abstract. In the paper, weak solvability of the first boundary value problem is proved for a class of parabolic equations with discontinuous coefficients and given in parabolic type domains in Sobolev's weight spaces. The coefficients of these equation bear discontinuity at the vertex of P- domain. At the vertex P- domain touches the characteristics of the equation.

Key Words and Phrases: boundary value problem, weak solvability, parabolic operator. 2010 Mathematics Subject Classifications: 35K05, 35K08, 34B05

1. Introduction

Let E_n and R_{n+1} be -n- dimensional and (n+1) dimensional Euclidean spaces of the points $x = (x_1, \ldots, x_n)$ and $(x, t) = (x_1, \ldots, x_n, t)$ respectively. D be a bounded domain E_n with a boundary ∂D , $0 \in D$, $R_{n+1}^- = R_{n+1} \cap \{(x, t) : t < 0\}$.

The domain $Q \subset R_{n+1}^-$ is said to be a paraboloid type domain (or P-domain) if its cross section with each hyperplane $t = \tau (\tau < 0)$ has the form:

$$\left\{x: \frac{x}{2\sqrt{-\tau}} \in D\right\}.$$

The domain D is called a foot of the P- domain Q. Let further $Q_T = Q \cap \{(x,t) : -T < t < 0\}, \quad S_T = \partial Q \cap \{(x,t) : T < t < 0\},$ $D_T = Q \cap \{(x,t) : t = -T\}, \quad \Gamma(Q_T)$ be a parabolic boundary of the domain Q_T . Consider in Q_T the following operator

$$L = \Delta + \lambda \sum_{i,j=1}^{n} \frac{x_i x_j}{4(-t)} \cdot \frac{\partial^2}{\partial x_i \partial x_j} - \frac{\partial U}{\partial t},$$

where Δ is the Laplace operator and the number parameter λ satisfies the condition

$$\frac{1}{d^2} < \lambda < \infty. \tag{1}$$

http://www.cjamee.org

42

© 2013 CJAMEE All rights reserved.

Here $d = \sup_{y \in D} |y|$. It is easy to see that subject to condition (1) the operator L uniformly parabolic in the domain Q_T . By analogy with the elliptic case, we call the operator L the Gilbarg-Serrin parabolic operator.

Let us agree in the following denotation u_i and u_{ij} are the derivatives of $\frac{\partial u}{\partial x_i}$ and $\frac{\partial^2 u}{\partial x_i \partial x_j}$, respectively,

$$u_{xx} = (u_{ij}), \qquad u_x^2 = \sum_{i=1}^n u_i^2, \quad u_{xx}^2 = \sum_{i,j=1}^n u_{ij}^2; \quad i, j = \overline{1, n}.$$

Let the number parameter γ satisfy the condition

$$\gamma \epsilon \left(\frac{n^2 \left(\lambda - \frac{1}{d^2} \right) + 2\lambda n}{2}, \infty \right).$$
⁽²⁾

 $A_0^{\infty}(Q_T)$ be a space of infinitely differentiable and finite in Q_T functions for which the following integral is finite $\int_{Q_T} (-t)^{\gamma} u^2 dx dt$, $L_{2,\gamma}(Q_T)$ be Banach space of measurable functions u(x,t) given on, Q_T with finite norm

$$\|u\|_{L_{2,\gamma}(Q_T)} = \left(\int_{Q_T} (-t)^2 u^2 dx dt\right)^{\frac{1}{2}},$$

 $\overset{0}{W_{2,\gamma}}^{1,0}(Q_T)$ and $\overset{0}{W_{2,\gamma}}^{1,1}(Q_T)$ be Banach spaces of measurable functions u(x,t) given on Q_T with finite norms

$$\begin{split} \|u\|_{W^{1,0}_{2,\gamma}(Q_T)} &= \left(\int_{Q_T} (-t)^{\gamma} (u^2 + u_x^2) dx dt\right)^{\frac{1}{2}}, \\ \|u\|_{W^{1,1}_{2,\gamma}(Q_T)} &= \left(\int_{Q_T} (-t)^{\gamma} (u^2 + u_x^2 + u_t^2) dx dt\right)^{\frac{1}{2}}, \end{split}$$

respectively.

 $W_{2,\gamma}^{(0)}(Q_T)$ and $W_{2,\gamma}^{(1)}(Q_T)$ be subspaces of $W_{2,\gamma}^{(1)}(Q_T)$ and $W_{2,\gamma}^{(1)}(Q_T)$, respectively, in which $A_0^{\infty}(Q_T)$ is a dense set.

In the domain Q_T consider the first boundary value problem

$$Lu = \Delta u + \lambda \sum_{i,j=1}^{n} \frac{x_i x_j}{4(-t)} \cdot \frac{\partial^2 u}{\partial x_i \partial x_j} - \frac{\partial u}{\partial t} = f + \sum_{k=1}^{n} \frac{\partial f^k}{\partial x^k}$$
(3)

$$u|_{\Gamma(Q_T)} = 0, \tag{4}$$

where $f \epsilon L_{2,\gamma}(Q_T)$, $f^k \epsilon L_{2,\gamma}(Q_T)$; $k = \overline{1, n}$.

Therewith, it is assumed that with regard to number parameters λ and γ , conditions (1) and (2) are fulfilled. At first give definition of the weak solution of the first boundary value problem (3)-(4).

The function $u(x,t)\epsilon W_{2,\gamma}^{1,0}(Q_T)$ is said to be a weak solution of equation (3) in the domain Q_T if for any function $v(x,t)\epsilon W_{2,\gamma}^{0,1,1}(Q_T)$ the following integral identity is fulfilled.

$$B_{Q_{T}}(u,v) = \int_{Q_{T}} (-t)^{\gamma} u \, v_{t} dx dt - \int_{Q_{T}} (-t)^{\gamma} \sum_{i,j=1}^{n} \left(\delta_{ij} + \lambda \frac{x_{i} x_{j}}{4(-t)} \right) v_{i} u_{j} dx dt + + \lambda \, (n+1) \int_{Q_{T}} (-t)^{\gamma} \sum_{i=1}^{n} \frac{x_{i}}{4(-t)} u v_{i} dx dt + \frac{\lambda n \, (n+1)}{4} \int_{Q_{T}} (-t)^{\gamma} u v dx dt - - \gamma \int_{Q_{T}} (-t)^{\gamma} u v dx dt = \int_{Q_{T}} (-t)^{\gamma} f v dx dt - \int_{Q_{T}} (-t)^{\gamma} \sum_{k=1}^{n} f^{k} v_{k} dx dt,$$
(5)

where δ_{ij} is the Kronecker symbol.

The function $u(x,t)\epsilon W_{2,\gamma}^{(0)}(Q_T)$ being a weak solution of equation (3) in Q_T is called a weak solution of boundary value problem (3)-(5). Now we show the relation between equation (3) and integral identity (5). At first we represent the Hilbarg-Serrin parabolic operator in the form of a divergent operator with unbounded minor coefficients. We have

$$Lu = \Delta u + \lambda \sum_{i,j=1}^{n} \left(\frac{x_i x_j}{4(-t)} u_i \right)_j - \lambda (n+1) \sum_{i=1}^{n} \frac{x_i}{4(-t)} u_i - u_t .$$

Consider the domain $Q_{T,\delta} = Q_T \setminus \overline{Q_\delta}$ multiply the both parts of equation (3) by the function $v(x,t) \epsilon A_0^{\infty}(Q_T)$ and integrate the obtained equality with respect to $Q_{T,\delta}$. We get

$$\int_{Q_{T,\delta}} (-t)^{\gamma} v \ \Delta u \ dx \ dt +$$

$$+\lambda \int_{Q_{T,\delta}} (-t)^{\gamma} \sum_{i,j=1}^{n} \left(\frac{x_{??} x_j}{4 (-t)} u_i \right)_j v \, dx \, dt - \lambda (n+1) \int_{Q_{T,\delta}} (-t)^{\gamma} \sum_{i=1}^{n} \frac{x_j}{4 (-t)} u_j v \, dx \, dt - \int_{Q_{T,\delta}} (-t)^{\gamma} u_t v \, dx \, dt = \int_{Q_{T,\delta}} (-t)^{\gamma} f \cdot v \, dx \, dt + \int_{Q_{T,\delta}} (-t)^{\gamma} \sum_{i=1}^{n} \frac{\partial f^i}{\partial x_i} v \, dx \, dt.$$
(6)

By Ostrogradskii's formula

$$\int_{Q_{T,\delta}} (-t)^{\gamma} v \,\Delta u \,dx \,dt + \lambda \int_{Q_{T,\delta}} (-t)^{\gamma} \sum_{i,j=1}^{n} \left(\frac{x_i x_j}{4 \, (-t)} u_i \right)_j v \,dx \,dt =$$
$$= -\int_{Q_{T,\delta}} (-t)^{\gamma} \sum_{i,j=1}^{n} \left(\delta_{ij} + \lambda \frac{x_i x_j}{4 \, (-t)} \right) u_i v_j \,dx \,dt. \tag{7}$$

In what follows we have

$$-\lambda (n+1) \int_{Q_{T,\delta}} (-t)^{\gamma} \sum_{i=1}^{n} \frac{x_i}{4(-t)} u_i v \, dx \, dt = \frac{\lambda n (n+1)}{4} \int_{Q_{T,\delta}} (-t)^{\gamma-1} u v \, dx \, dt + \lambda (n+1) \int_{Q_{T,\delta}} (-t)^{\gamma} \sum_{i=1}^{n} \frac{x_i}{4(-t)} v_i u \, dx \, dt$$
(8)

Furthermore

$$\int_{Q_{T,\delta}} (-t)^{\gamma} f \cdot v \, dx \, dt + \int_{Q_{T,\delta}} (-t)^{\gamma} \sum_{i=1}^{n} \frac{\partial f^{i}}{\partial x_{i}} v \, dx \, dt =$$
$$= \int_{Q_{T,\delta}} (-t)^{\gamma} f \, v \, dx \, dt - \int_{Q_{T,\delta}} (-t)^{\gamma} \sum_{i=1}^{n} f^{i} \, v_{i} \, dx \, dt \tag{9}$$

Let $\Pi_R = \{x : |x_i| < R, i = \overline{1, n}\}, \quad K_{\delta} = \Pi_R \times (-T, -\delta).$ For simplicity we will consider that we can continue the function u(x, t) in $K_{\delta} \setminus Q_{T,\delta}$ so that the obtained continuation $\tilde{u}(x, t)$ be the element of the space $W_{2,\gamma}^{1,0}(K_{\delta})$. We continue the function v(x, t) by a zero to $K_{\delta} \setminus Q_{T,\delta}$ and denote the obtained continuation again by v(x, t). We have

$$J_{\delta} = -\int_{K_{\delta}} (-t)^{\gamma} \tilde{u}_{t} v \, dx \, dt = -\delta \int_{\Pi_{R}} \widetilde{D} (x, -\delta) \, v \, (x, -\delta) \, dx + \\ + \int_{K_{\delta}} (-t)^{\gamma} v_{t} \tilde{u} \, dx \, dt - \gamma \int_{K_{\delta}} (-t)^{\gamma-1} \widetilde{u} \, v \, dx \, dt = -\delta^{\gamma} \int_{K_{\delta}} u (x, -\delta) \, v \, (x, -\delta) \, dx + \\ + \int_{Q_{T,\delta}} (-t)^{\gamma} v_{t} u dx dt - \gamma \int_{Q_{T,\delta}} (-t)^{\gamma-1} u v dx dt.$$

Hence it follows that

$$\lim_{\delta \to 0+} J_{\delta} = \int_{Q_T} (-t)^{\gamma} v_t \ u \ dx \ dt - \gamma \int_{Q_T} (-t)^{\gamma-1} u \ v \ dx \ dt.$$
(10)

Now, taking into account (7)-(10) in (6), and tending δ to zero we arrive at integral identity (5).

Theorem 1. If with respect to number parameters λ and γ conditions (1) and (2) are fulfilled, then the first boundary value problem (3)-(4) is uniquely weakly solvable in the $_{0}^{0,0}$ space $W_{2,\gamma}(Q_T)$ for any $f(x,t) \in L_{2,\gamma}(Q_T)$ and $f^k(x,t) \in L_{2,\gamma}(Q_T)$; $k = \overline{1,n}$.

Proof. At first prove the existence of the solution. To this end we consider the extending sequence of domains $\{D_m\}$, m = 1, 2, ...; approximating from within the domain D, i.e. $\overline{D_m} \subset D_{m+1}$, $\overline{D_m} \subset D$, $\lim_{m\to\infty} D_m = D$. Therewith we choose D_m so that for any natural $m \quad \partial D_m \in C^2$. Let further Q^m be a P-domain whose foot is the domain D_m ,

$$Q_T^m = Q^m \cap \{(x,t) : t > -T\}, \quad Q_{T,\delta}^m = Q_T^m \setminus \overline{Q}_{\delta}^m, \quad \delta \in (0,T).$$

Denote by f^h and $f^{k,h}$ the Friedrichs averaged functions, respectively, $k = \overline{1, n}$ with a parameter h > 0. Consider for h > 0 and natural m the family of the first boundary value problems

$$\Delta u^{m,h} + \lambda \sum_{i,j=1}^{n} \left(\frac{x_i x_j}{4(-t)} u_i^{m,h} \right)_j - \lambda (n+1) \sum_{i=1}^{n} \frac{x_i}{4(-t)} u_i^{m,h} - u_t^{m,h} =$$
$$= f^h + \sum_{k=1}^{n} \frac{\partial f^{k,h}}{\partial x_k} ; \quad (x,t) \in Q^m_{T,\delta},$$
(11)

$$u^{m,h}\Big|_{\Gamma\left(Q^m_{T,\delta}\right)} = 0.$$
⁽¹²⁾

As for any natural m and positive h and δ the coefficients and the right side of equation (11) are infinitely differentiable in $\overline{Q}_{T,\delta}^m$ functions, problem (11)-(12) has a unique classic solution $u^{m,h}(x,t)$. Indeed, $u^{m,h}(x,t)$ depends on δ as well, but for brevity of notation we write $u^{m,h}(x,t)$ instead of $u_{\delta}^{m,h}(x,t)$. Multiply the both sides of equation (11) by the function $(-t)^{\gamma}u^{m,h}(x,t)$ and integrate the obtained equality with respect to the domain $\begin{array}{c} Q^m_{T,\delta}.\\ \text{We get} \end{array}$

$$\int_{Q_{T,\delta}^{m}} (-t)^{\gamma} \Delta u^{m,h} \cdot u^{m,h} dx \, dt + \int_{Q_{T,\delta}^{m}} (-t)^{\gamma} u^{m,h} \sum_{i,j=1}^{n} \frac{\partial}{\partial x_{i}} \left(\frac{x_{i}x_{j}}{4(-t)} u_{j}^{m,h} \right) dx \, dt - \\
-\lambda (n+1) \int_{Q_{T,\delta}^{m}} (-t)^{\gamma} \sum_{i=1}^{n} \frac{x_{i}}{4(-t)} u_{i}^{m,h} u^{m,h} dx \, dt - \int_{Q_{T,\delta}^{m}} (-t)^{\gamma} u_{i}^{m,h} u_{t}^{m,h} dx dt = \\
= \int_{Q_{T,\delta}^{m}} (-t)^{\gamma} f^{h} \cdot u^{m,h} dx \, dt + \int_{Q_{T,\delta}^{m}} (-t)^{\gamma} \sum_{k=1}^{n} \frac{\partial f^{k,h}}{\partial x_{k}} u^{m,h} dx \, dt. \tag{13}$$

Further we have

$$\int_{Q_{T,\delta}^{m}} (-t)^{\gamma} \Delta u^{m,h} \cdot u^{m,h} dx \, dt + \lambda \int_{Q_{T,\delta}^{m}} (-t)^{\gamma} u^{m,h} \sum_{i,j=1}^{n} \frac{\partial}{\partial x_{i}} \left(\frac{x_{i}x_{j}}{4(-t)} u_{j}^{m,h} \right) dx \, dt = \\ = -\int_{Q_{T,\delta}^{m}} (-t)^{\gamma} \left(u_{x}^{m,h} \right)^{2} dx \, dt - \lambda \int_{Q_{T,\delta}^{m}} (-t)^{\gamma} \sum_{i,j=1}^{n} \frac{x_{i}x_{j}}{4(-t)} u_{i}^{m,h} u_{j}^{m,h} dx dt.$$
(14)

Furthermore

$$-\lambda (n+1) \int_{Q_{T,\delta}^m} (-t)^{\gamma} \sum_{i=1}^n \frac{x_i}{4(-t)} u_i u \, dx dt = \frac{\lambda (n+1) \cdot n}{2} \int_{Q_{T,\delta}^m} (-t)^{\gamma} \frac{u^2}{4(-t)} dx \, dt, \quad (15)$$

$$\int_{Q_{T,\delta}^m} (-t)^{\gamma} \sum_{k=1}^n \frac{\partial f^{k,h}}{\partial x_k} u^{m,k} dx \ dt = -\int_{Q_{T,\delta}^m} (-t)^{\gamma} \sum_{k=1}^n f^{k,h} u_k^{m,h} dx \ dt.$$
(16)

Weak Solvability of the First Boundary Value Problem for a Class of Parabolic Equations 47

Finally, by means of arguments similar to ones that were used by deriving integral identity (5), we get

$$-\int_{Q_{T,\delta}^{m}} (-t)^{\gamma} u \ u_{t} dx dt = -\frac{\gamma}{2} \int_{Q_{T,\delta}^{m}} (-t)^{\gamma} u^{2} dx dt + i_{1} \left(\delta\right), \tag{17}$$

where $\lim_{\delta \to \infty} i_1(\delta) = 0$. Taking into account (13)-(16) in (12) and tending δ to zero, we conclude

$$\int_{Q_{T,\delta}^{m}} (-t)^{\gamma} \left(\left(u_{x}^{m,h} \right)^{2} + \lambda \sum_{i,j=1}^{n} \frac{x_{i}x_{j}}{4(-t)} u_{i}^{m,h} \cdot u_{j}^{m,h} \right) dx dt - \frac{\lambda (n+1) - 4\gamma}{2} \cdot \int_{Q_{T}^{m}} (-t)^{\gamma} \frac{(u^{m,h})^{2}}{4(-t)} dx dt = \int_{Q_{T}^{m}} (-t)^{\gamma} \sum_{k=1}^{n} f^{k} u_{k}^{m,h} dx dt - \int_{Q_{T}^{m}} (-t)^{\gamma} f u^{m,h} dx dt.$$
(18)

Here $u^{m,h}(x,t) = \lim_{\delta \to 0+} u_{\delta}^{m,h}(x,t)$. The existence of the pointwise limit is proved in the same as in [1]. If now $\frac{\lambda n(n+1)-4\gamma}{2} \leq 0$, i.e. $\gamma \geq \frac{\lambda n(n+1)}{2}$, then from (17) it follows

$$\int_{Q_T^m} (-t)^{\gamma} \left(\left(u_x^{m,h} \right)^2 + \lambda \sum_{i,j=1}^n \frac{x_i x_j}{4 (-t)} u_i^{m,h} \cdot u_j^{m,h} \right) dx \, dt \leq \\
\leq \int_{Q_T^m} (-t)^{\gamma} \sum_{k=1}^n f^{k,h} u_k^{m,h} dx \, dt - \int_{Q_T^m} (-t)^{\gamma} f^h u^{m,h} dx \, dt. \tag{19}$$

Note that for ≥ 0 , $\lambda \sum_{i,j=1}^{n} \frac{x_i x_j}{4(-t)} u_i^{m,h} \cdot u_j^{m,h} \geq 0$. But if $-\frac{1}{d^2} < \lambda < 0$, then

$$\lambda \sum_{i,j=1}^{n} \frac{x_i x_j}{4(-t)} u_i^{m,h} \cdot u_j^{m,h} \ge \lambda \ d^2 \left(u_x^{m,h} \right)^2.$$

Thus, if $\gamma \geq \frac{\lambda n(n+1)}{2}$, then from (18) we get

$$\int_{Q_T^m} (-t)^{\gamma} \left(u_x^{m,h} \right)^2 dx \, dt \leq C_1 \left(\lambda, n, d, \gamma \right).$$

$$\left(\int_{Q_T^m} (-t)^{\gamma} \sum_{k=1}^n f^{k,h} u_k^{m,h} dx \, dt - \int_{Q_T^m} (-t)^{\gamma} f^h u^{m,h} dx \, dt \right). \tag{20}$$

Now consider the case

$$\gamma \in \left(\frac{n^2 \left(\lambda - \frac{1}{d^2}\right) + 2\lambda n}{8}, \frac{\lambda n \left(n+1\right)}{2}\right).$$
(21)

According to inequality (17)

$$\frac{\lambda n (n+1) - 4\gamma}{2} \int_{Q_T^m} (-t)^{\gamma} \frac{(u^{m,h})^2}{4(-t)} dx dt \leq \\ \leq \frac{2 \lambda n (n+1) - 8\gamma}{n^2} \int_{Q_T^m} (-t)^{\gamma} \sum_{i,j=1}^n \frac{x_i x_j}{4(-t)} u_i^{m,h} \cdot u_j^{m,h} dx dt.$$
(22)

But on the other hand, from (20) it follows that there exists $\mu \in (0, 1)$ for which

$$\frac{2\ \lambda\ n\left(n+1\right)-8\gamma}{n^2} < \frac{1}{d^2} + \lambda - \frac{\mu}{d^2}.$$

So, from (18) and (21) we conclude

$$\int_{Q_T^m} (-t)^{\gamma} \left(\left(u_x^{m,h} \right)^2 + \frac{\mu - 1}{d^2} \sum_{i,j=1}^n \frac{x_i x_j}{4 (-t)} u_i^{m,h} \cdot u_j^{m,h} dx \, dt \right) \leq \\
\leq \int_{Q_T^m} (-t)^{\gamma} \left(\sum_{i=1}^n f^{i,h} u_i^{m,h} - f^h u^{m,h} \right) dx \, dt.$$
(23)

As $\mu < 1$, then

$$\frac{\mu - 1}{d^2} \sum_{i,j=1}^n \frac{x_i x_j}{4(-t)} u_i^{m,h} \cdot u_j^{m,h} \ge (\mu - 1) \left(u_x^{m,h} \right)^2.$$
(24)

From (22)-(23) it follows that

$$\mu \int_{Q_T^m} (-t)^{\gamma} \left(u_x^{m,h} \right)^2 dx \ dt \le \int_{Q_T^m} (-t)^{\gamma} \left(\sum_{k=1}^n f^{k,h} u_k^{m,h} - f^h u^{m,h} \right) dx \ dt.$$

The last inequality and estimation (19) allows to conclude that for $\gamma \in \left(\frac{n^2\left(\lambda - \frac{1}{d^2}\right) + 2\lambda n}{8}, \infty\right)$ the following inequality is valid

$$\int_{Q_T^m} (-t)^{\gamma} \left(u_x^{m,h} \right)^2 dx \ dt \le C_2 \left(\lambda, n, d, \gamma \right) \int_{Q_T^m} (-t)^{\gamma} \left(\sum_{k=1}^n f^{k,h} u_k^{m,h} - f^h u^{m,h} \right) dx \ dt.$$
(25)

Weak Solvability of the First Boundary Value Problem for a Class of Parabolic Equations 49

According to Friedrich's inequality we get

$$\int_{Q_T^m} (-t)^{\gamma} \left(u^{m,h} \right)^2 dx \ dt \le C_3 \left(\lambda, n, d, \gamma \right) \int_{Q_T^m} (-t)^{\gamma} \left(\sum_{k=1}^n f^{k,h} u_k^{m,h} - f^h u^{m,h} \right) dx \ dt.$$
(26)

Thus, from (24)-(25) we conclude

$$\int_{Q_T^m} (-t)^{\gamma} \left(\left(u^{m,h} \right)^2 + \left(u_x^{m,h} \right)^2 \right) dx \ dt \le leC_4 \left(\lambda, n, d, \gamma \right) \int_{Q_T^m} (-t)^{\gamma} \left(\sum_{k=1}^n f^{k,h} u_k^{m,h} - f^h u^{m,h} \right) dx \ dt.$$

$$(27)$$

Further, for any $\varepsilon > 0$ we have

$$\int_{Q_T^m} (-t)^{\gamma} \left(\sum_{k=1}^n f^{k,h} u_k^{m,h} - f^h u^{m,h} \right) dx \, dt \leq \\ \leq \frac{\varepsilon}{2} \int_{Q_T^m} (-t)^{\gamma} \sum_{k=1}^n \left(u_k^{m,h} \right)^2 dx \, dt + \frac{1}{2\varepsilon} \int_{Q_T^m} (-t)^{\gamma} \sum_{k=1}^n \left(f^{k,h} \right)^2 dx \, dt + \\ + \frac{\varepsilon}{2} \int_{Q_T^m} (-t)^{\gamma} \left(u^{m,h} \right)^2 dx \, dt + \frac{1}{2\varepsilon} \int_{Q_T^m} (-t)^{\gamma} \left(u^{m,h} \right)^2 dx \, dt.$$

$$(28)$$

Now choosing $\varepsilon = \frac{1}{C_4}$ from (26)-(27) we get

$$\int_{Q_T^m} (-t)^{\gamma} \left(\left(u^{m,h} \right)^2 + \left(u_x^{m,h} \right)^2 \right) dx \ dt \le C_5 \left(\lambda, n, d, \gamma \right) \times \\ \times \left(\int_{Q_T^m} (-t)^{\gamma} \sum_{k=1}^n \left(f^{k,h} \right)^2 dx \ dt + \int_{Q_T^m} (-t)^{\gamma} \left(f^h \right)^2 dx \ dt \right).$$

$$\tag{29}$$

Without loss of generality, we can consider that for $f \neq 0$; $f^k \neq 0$; $k = \overline{1, n}$. Therefore from (28) it follows that for rather small h > 0

$$\left\| u^{m,k} \right\|_{W^{1,0}_{2,\gamma}(Q^m_T)} \le C_6\left(\lambda, n, d, \gamma\right) \left(\|f\|_{L_{2,\gamma}(Q_T)} + \sum_{k=1}^n \left\| f^k \right\|_{L_{2,\gamma}(Q_T)} \right).$$
(30)

Fix an arbitrary natural m. From inequality (29) it follows that the family of functions $\{u^{m,h}(x,t)\}$ is weakly compact (with respect to h) in the space $W_{2,\gamma}^{0,0}(Q_T^m)$. Thus, there exists such a sequence $h_l \to 0$ as $l \to \infty$ and the function $u^m(x,t) \in W_{2,\gamma}^{0,0}(Q_T^m)$ that the functional sequence $\{u^{m,h_l}(x,t)\}$ weakly converges to the function $u^m(x,t)$ in $W_{2,\gamma}^{0,0}(Q_T^m)$

as $l \to \infty$. This means that for any function $u^m(x,t) \in \overset{0}{W}^{1,0}_{2,\gamma}(Q^m_T)$ it holds the limit equality

$$\lim_{l \to \infty} B_{Q_T^m}\left(u^{m,h_l}, v\right) = B_{Q_T^m}\left(u^m, v\right) \,. \tag{31}$$

But on the other hand

$$B_{Q_T^m}\left(u^{m,h_l},v\right) = \int_{Q_T^m} (-t)^{\gamma} f^{h_l} v \, dx \, dt - \int_{Q_T^m} (-t)^{\gamma} \sum_{k=1}^n f^{k,h_l} v_k dx \, dt.$$
(32)

Furthermore

$$\lim_{l \to \infty} \left(\int_{Q_T^m} (-t)^{\gamma} f^{h_l} v \, dx \, dt - \int_{Q_T^m} (-t)^{\gamma} \sum_{k=1}^n f^{k,h_l} v_k \, dx \, dt \right) = \int_{Q_T^m} (-t)^{\gamma} f v \, dx \, dt - \int_{Q_T^m} (-t)^{\gamma} \sum_{k=1}^n f^k v_k dx \, dt.$$
(33)

From (30)-(32) we conclude that

$$B_{Q_T^m}(u^m, v) = \int_{Q_T^m} (-t)^{\gamma} f v \, dx \, dt - \int_{Q_T^m} (-t)^{\gamma} \sum_{k=1}^n f^k v_k dx \, dt$$

The last equality means that the function $u^{m}(x,t)$ is a weak solution of equation (3) in the domain Q_{T}^{m} . Furthermore, for the function $u^{m}(x,t)$ the following estimation is valid

$$\|u^{m}\|_{W^{1,0}_{2,\gamma}(Q_{T}^{m})} \leq C_{7}\left(\lambda, n, d, \gamma\right) \left(\|f\|_{L_{2,\gamma}(Q_{T})} + \sum_{k=1}^{n} \left\|f^{k}\right\|_{L_{2,\gamma}(Q_{T})}\right).$$
(34)

For any natural m we continue the function $u^m(x,t)$ by a zero in $Q_T \setminus Q_T^m$ and denote the obtained continuation again by $u^m(x,t)$. It is easy to see that $u^m(x,t) \in W_{2,\gamma}^{0,1,0}(Q_T)$. Therewith, according to (33) the following estimation is valid

$$\|u^{m}\|_{W^{1,0}_{2,\gamma}(Q_{T})} \leq C_{8} \left(\|f\|_{L_{2,\gamma}(Q_{T})} + \sum_{k=1}^{n} \left\|f^{k}\right\|_{L_{2,\gamma}(Q_{T})} \right).$$
(35)

From (34) it follows that the family of functions $\{u^m(x,t)\}, \ldots, m = 1, 2, \ldots$ is weakly compact in the space $W_{2,\gamma}(Q_T)$. Thus, there exists such a function $u(x,t) \in W_{2,\gamma}(Q_T^m)$ and sequence $m_r \to \infty$ as $r \to \infty$ that u(x,t) is a weak limit of $u^{m_r}(x,t)$ as $r \to \infty$ in $0^{1,0}$ $W_{2,\gamma}(Q_T^m)$. This means that for any function $v(x,t) \in W_{2,\gamma}$ the following limit equality is valid:

$$\lim_{r \to \infty} B_{Q_T} \left(u^{m_r}, v \right) = B_{Q_T} \left(u, v \right) \,.$$

Weak Solvability of the First Boundary Value Problem for a Class of Parabolic Equations

Moreover, using the above arguments, we can show that

$$B_{Q_T}(u,v) = \int_{Q_T^m} (-t)^{\gamma} f v \, dx \, dt - \int_{Q_T^m} (-t)^{\gamma} \sum_{k=1}^n f^k v \, dx \, dt.$$

From the last equality it follows that the function u(x,t) is a weak solution of the first boundary value problem (3)-(4). Furthermore, for the functions u(x,t) the following estimation is valid

$$\|u^{m}\|_{U_{2,\gamma}^{0,1,0}(Q_{T}^{m})} \leq C_{9}\left(\|f\|_{L_{2,\gamma}(Q_{T})} + \sum_{k=1}^{n} \left\|f^{k}\right\|_{L_{2,\gamma}(Q_{T})}\right).$$
(36)

Thereby the existence of the weak solution of the first boundary value problem (3)-(4) is proved. Now prove its uniqueness. It suffices to show that a homogeneous problem has only a trivial solution. Let u(x,t) be the solution of homogeneous problem (3)-(4), i.e. for $f \equiv 0$; $f^k \equiv 0$; $k = \overline{1, n}$. Fix an arbitrary $\delta \in (0, T)$ and consider the function $v(x,t) \in W_{2,\gamma}(Q_{T+\delta})$ vanishing for $t \leq T$ and $t \geq -\delta$. Let further $K = \prod_R \times (-T - \delta, 0)$, $\prod_R = \{x : |x_i| < R, i = \overline{1, n}\}$. Continue the function u(x,t) and v(x,t) by zero to $K \setminus Q_T$ and denote the obtained continuations again by u(x,t) and v(x,t), respectively. It is easy to see that $u(x,t) \in W_{2,\gamma}(K)$, while $v(x,t) \in W_{2,\gamma}(K)$, Denote for

$$h \in (0, \delta] \frac{1}{h} \int_{t-h}^{t} v(x, \tau) d\tau$$
 by $v_{\overline{h}}(x, \tau)$

and put into integral identity (5) instead of the function u(x,t) the function $v_{\overline{h}}(x,\tau)$ and get

$$B_K\left(u, v_{\overline{h}}\right) = 0. \tag{37}$$

Taking into account the equalities $(v_{\overline{h}})_t = (v_t)_{\overline{h}}$, $(v_{\overline{h}})_i = (v_i)_{\overline{h}}$ $i = \overline{1, n}$ and also

$$-\int_{K} (-t)^{\gamma} u \ (v_{t})_{\overline{h}} \ dx \ dt = -\int_{K} ((-t)^{\gamma} u)_{h} v_{t} dx \ dt = \int_{K} [((-t)^{\gamma} u)_{h}]_{t} v dx \ dt,$$

$$\int_{K} (-t)^{\gamma} \sum_{i,j=1}^{n} \left(\delta_{ij} + \lambda \frac{x_{i} x_{j}}{4 \ (-t)} \right) u_{i} u_{i} \ (v_{j})_{\overline{h}} \ dx \ dt = \int_{K} (-t)^{\gamma} \sum_{i,j=1}^{n} \left(\delta_{ij} + \lambda \frac{x_{i} x_{j}}{4 \ (-t)} u_{i} \right) v_{j} \ dx \ dt,$$

where $u_h(x,t) = \frac{1}{h} \int_t^{t+h} u(x,\tau) d\tau$, assuming $v(x,t) = u_h(x,t)$ tending h to zero, from (36) we get

$$\frac{\lambda n (n+1) - 4\gamma}{8} \int_{Q_{T,\delta}} (-t)^{\gamma - 1} u^2 \, dx \, dt - \int_{Q_{T,\delta}} (-t)^{\gamma} \sum_{i,j=1}^n \left(\delta_{ij} + \lambda \frac{x_i x_j}{4 (-t)}\right) u_i u_j \, dx \, dt = 0.$$
(38)

Now, behaving as in deriving estimation (24), we get that if with respect to number parameters λ and γ conditions (1) and (2) are fulfilled, then $\int_{Q_{T,\delta}} (-t)^{\gamma} u_x^2 dx dt = 0$. The last equality yields

$$\int_{Q_{T,\delta}} (-t)^{\gamma-1} u^2 \ dx \ dt = 0$$

With regard to arbitrariness of δ we conclude

$$\int_{Q_T} (-t)^{\gamma - 1} u^2 \, dx \, dt = 0.$$

Hence it follows that u(x,t) = 0 almost everywhere in Q_T .

In fact in the course of proof we established the estimation of the weak solution of the first boundary value problem (3)-(4). We formulate this statement in the form of a separate theorem.

Theorem 2. If the conditions of the previous theorem are fulfilled then for the weak solution of the first boundary value problem (3)-(4), estimation (35) is valid.

References

- O.A. Ladyzhenskaya, V.A. Solonnikov, Uraltseva N.N. linear and quaslinear equations of parabolic type, Moscow, Nauka, 1967.
- [2] Yu.A. Alkhutov, Behavior of solutions of second order parabolic equations in noneylindrical domains, Doklady RAN, 345(5), 1995, 583-585.
- [3] Yu.A. Alkhutov, I.T. Mamedov, Some properties of the solution of the first boundary value problem for second order parabolic equations with discontinuous coefficients, Dokl. AN SSSR, 284(1), 1985, 477-500.
- [4] V.A. Rukavishnikov, U. Rukavishnikova, On the isomorphic mapping of weighted spaces by an elliptic operator with degeneration on the domain boundary, Diff. Urav., 50(3), 2014, 345-351.
- [5] A.B. Kostin, *Counterexamples in inverse problem for parabolic, elliptic, hyperbolic equations*, Zhurnal matematiki I matematicheskoi fiziki, **54(5)**, 2014, 797-810.
- [6] N.J. Jafarov, First boundary value problem for class of parabolic equations with discontinuous coefficients, the basic coercive estimation in Sobolev weight space, Trans. of NAS of Azerb. Ser. of phys. – tech. and math. sci., XXXV(1), 2015, 51-58.

Nazim J.Jafarov Institute of Mathematics and Mechanics of NAS of Azerbaijan, Baku, Azerbaijan E-mail: nazim.jafarov.math@gmail.com

Received 23 August 2018 Accepted 25 September 2018